

UNIVERSE+ Online Seminar

Agnese Bissi

“AdS Amplitudes from CFT”

Part 1

universe+ is a cooperation of

MAX PLANCK INSTITUTE
FOR PHYSICS



IAS

INSTITUTE FOR
ADVANCED STUDY

MAX PLANCK INSTITUTE
FOR MATHEMATICS IN THE SCIENCES



UNIVERSITY OF AMSTERDAM



Funded by
the European Union



European Research Council
Established by the European Commission

Funded by the European Union (ERC, UNIVERSE PLUS, 101118787). Views and opinions expressed are however those of the author(s) only and do not necessarily reflect those of the European Union or the European Research Council Executive Agency. Neither the European Union nor the granting authority can be held responsible for them.

AdS LOOPS AMPLITUDES FROM CFT

Agnese Bissi
(ICTP & UPPSALA UNIVERSITY)

BASED ON:

WORK IN PROGRESS W/ G. FARDELLI & M.R. KHANSARI
& ALSO DIFFERENT COLLABORATIONS W/ G. FARDELLI,
A. GEORGOUDIS, A. MANENTI

IN THE LAST YEARS THERE HAS BEEN AN INCREASING
INTEREST IN UNDERSTANDING HOW TO CLASSIFY AND
DEFINE CONFORMAL FIELD THEORIES IN $d \geq 3$.

IN THE LAST YEARS THERE HAS BEEN AN INCREASING
INTEREST IN UNDERSTANDING HOW TO CLASSIFY AND
DEFINE CONFORMAL FIELD THEORIES IN $d \geq 3$.

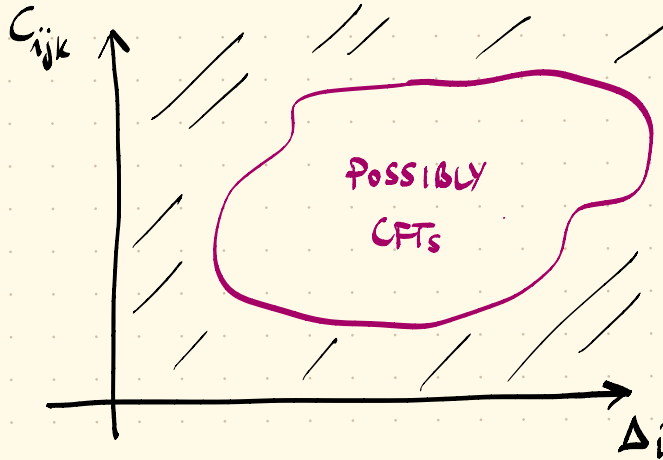
ONE IDEA: CONFORMAL BOOTSTRAP

CLASSIFY CFTs USING SYMMETRIES
AND CONSISTENCY CONDITIONS

CLASSIFY CFTs USING SYMMETRIES
AND CONSISTENCY CONDITIONS

CLASSIFY CFTs USING SYMMETRIES
AND CONSISTENCY CONDITIONS

SPECIFIED BY $\{\Delta_i, C_{ijk}\}$
→ OPE COEFFICIENTS
↳ CONFORMAL DIMENSIONS



/// → NO CONSISTENT CFTs

THIS APPROACH IS PARTICULARLY INTERESTING AND SUCCESSFULL

IN THE CONTEXT OF THE ADS/CFT CORRESPONDENCE.



HOW TO IDENTIFY
HOLOGRAPHIC CFTs
USING BOOTSTRAP TECHNIQUES



CONNECTIONS
WITH ADS
SIDE & CONSEQUENCES
FOR QUANTUM GRAVITY

AdS/CFT CORRESPONDENCE

AdS SCATTERING AMPLITUDES \leftrightarrow CFT CORRELATORS

LOOP EXPANSION
(IN G_N)



$\frac{1}{N}$ EXPANSION



c^{-1} EXPANSION



CENTRAL CHARGE OF
THE CFT

AdS/CFT CORRESPONDENCE

AdS SCATTERING AMPLITUDES \leftrightarrow CFT CORRELATORS

LOOP EXPANSION
(IN G_N)



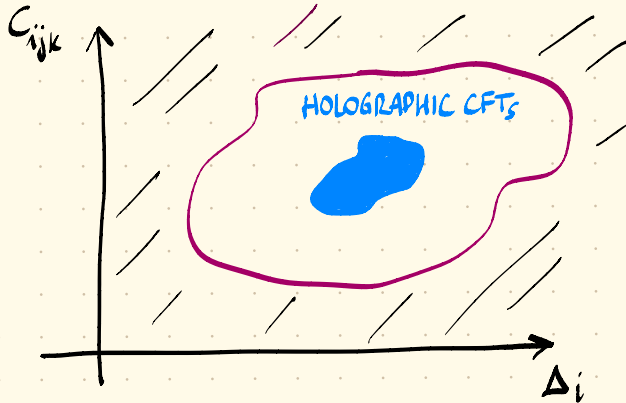
$\frac{1}{N}$ EXPANSION



C^{-1} EXPANSION



CENTRAL CHARGE OF
THE CFT



CONFORMAL FIELD THEORIES

- LOCAL OPERATORS

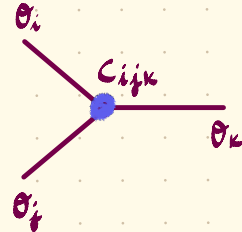
$\mathcal{O}_{\Delta, \ell}$
↑
DIMENSION
 ℓ SPIN



- THREE POINT FUNCTIONS

OPE COEFFICIENTS

C_{ijk}



↳ $\mathcal{O}_i \times \mathcal{O}_j \sim \boxed{C_{ijk}} \mathcal{O}_k + \dots$

WITH THIS PIECE OF INFORMATION, IT IS POSSIBLE TO
CONSTRUCT ALL HIGHER POINT FUNCTIONS,

- FOR INSTANCE 4 POINT FUNCTIONS

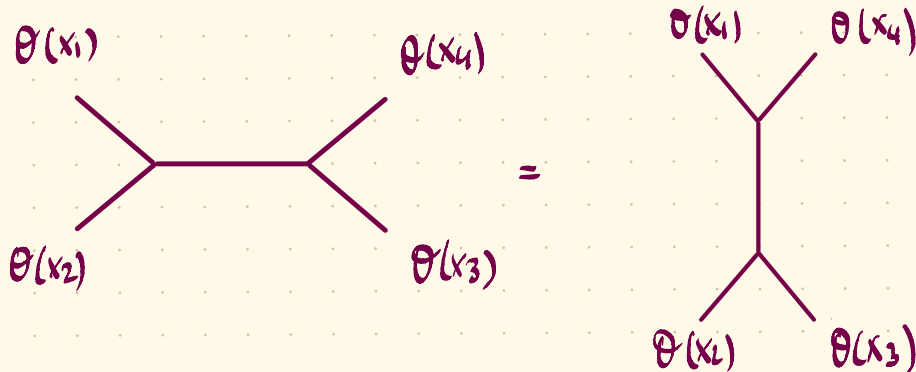
$$\text{Tree-level diagram} = \sum_j \text{Diagram with two vertices } c_{12j} \text{ and } c_{34j}$$

- THE SAME FOR HIGHER POINT FUNCTIONS

CONFORMAL BOOTSTRAP PHILOSOPHY :

- USE UNITARITY OF THE THEORY $\Delta \geq \Delta^u(d, l)$
 $c^2 \geq 0$

- USE ASSOCIATIVITY OF THE OPERATOR PRODUCT EXPANSION



- IF THERE ARE EXTRA SYMMETRIES, USE ALSO THEM (SUSY, $\mathcal{O}(N)$...)

AdS

CFT

picture

FIELDS ϕ_i

SINGLE TRACE OPERATORS \mathcal{O}_i

GRAVITON $g_{\mu\nu}$

STRESS TENSOR $T_{\mu\nu}$

MULTI PARTICLE STATES
(COMPOSITES)

HIGHER TRACES $[\mathcal{O}_i, \mathcal{O}_j, \dots, \mathcal{O}_k]$

MASS $m_i^2 = \Delta_i(\Delta_i - d)$

CONFORMAL DIMENSION Δ_i

AIM OF THIS TALK: • UNDERSTAND HOW TO USE
THE CFT DESCRIPTION
TO CONSTRAIN / CONSTRUCT
THE ADS AMPLITUDES

• FOCUS ON LARGE N LIMIT



SMALL G_N EXPANSION


LARGE N EXPANSION

$$\langle \theta(x_1) \theta(x_2) \theta(x_3) \theta(x_4) \rangle = \frac{G_f(u, v)}{\begin{matrix} 2\Delta\theta & 2\Delta\theta \\ x_{12} & x_{34} \end{matrix}}$$

LARGE N EXPANSION

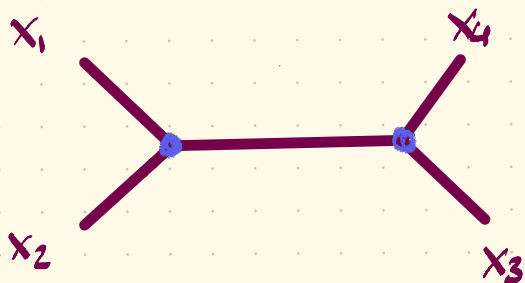
$$\langle \theta(x_1) \theta(x_2) \theta(x_3) \theta(x_4) \rangle = \frac{G(u, v)}{x_{12}^{2\Delta_\theta} x_{34}^{2\Delta_\theta}}$$

$$G(u, v) = \sum_{\Delta, e} a_{\Delta, e} \underbrace{u^{\frac{\Delta-e}{2}} \mathcal{G}_{\Delta, e}(u, v)}_{\text{CONFORMAL BLOCKS}} \quad d=4$$

$$G(u, v) = \sum_{\Delta, e} \text{Diagram}$$


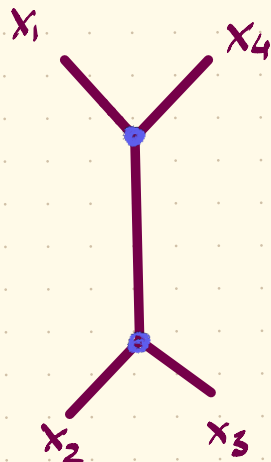
LARGE N EXPANSION

$$\langle \theta(x_1) \theta(x_2) \theta(x_3) \theta(x_4) \rangle = \frac{g(u, v)}{\frac{2\Delta\theta}{x_{12}} \frac{2\Delta\theta}{x_{34}}}$$



$$\frac{g(u, v)}{\frac{2\Delta\theta}{x_{12}} \frac{2\Delta\theta}{x_{34}}}$$

=



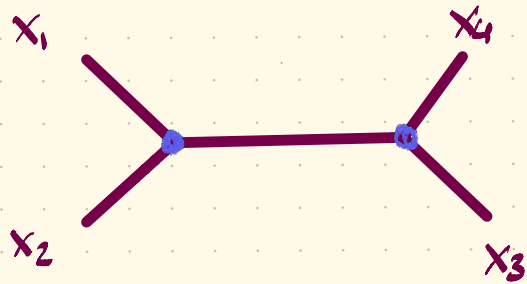
$$\frac{g(v, u)}{\frac{2\Delta\theta}{x_{14}} \frac{2\Delta\theta}{x_{34}}}$$

$$u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2}$$

$$v = \frac{x_{14}^2 x_{23}^2}{x_{13}^2 x_{24}^2}$$

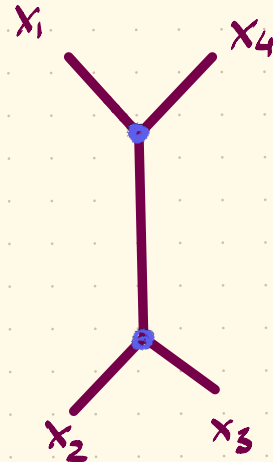
LARGE N EXPANSION

$$\langle \theta(x_1) \theta(x_2) \theta(x_3) \theta(x_4) \rangle = \frac{g(u, v)}{\frac{2\Delta\theta}{x_{12}} \frac{2\Delta\theta}{x_{34}}}$$



$g(u, v)$

$=$



$g(v, u)$

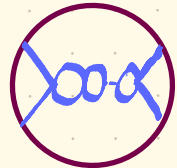
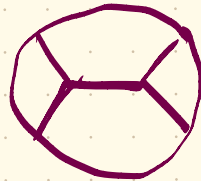
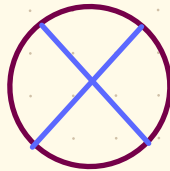
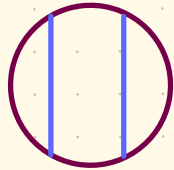
$\left(\frac{u}{v}\right)^{\Delta\theta}$

$$u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2}$$

$$v = \frac{x_{14}^2 x_{23}^2}{x_{13}^2 x_{24}^2}$$

- EXPAND THE CORRELATOR AND OPE DATA IN THE LARGE N LIMIT

$$G(u, v) = G^{(0)}(u, v) + \frac{1}{N^2} G^{(1)}(u, v) + \dots + \frac{1}{N^{2k}} G^{(k)}(u, v)$$



$$\Delta = \Delta^{(0)} + \frac{1}{N^2} \delta^{(1)} + \dots + \frac{1}{N^{2k}} \delta^{(k)}$$

$$a_{\Delta, e} = a^{(0)} + \frac{1}{N^2} a^{(1)} + \dots + \frac{1}{N^{2k}} a^{(k)}$$

∴ higher trace + new operators

- DEPENDING AT WHICH ORDER WE ARE, THE INTERMEDIATE OPERATORS ARE DIFFERENT

- DEPENDING AT WHICH ORDER WE ARE, THE INTERMEDIATE OPERATORS ARE DIFFERENT

- GENERICALLY:

→ AT LEADING ORDER N^0 ONLY DOUBLE TRACE OPERATORS → $[\theta\theta]_{n,l} \equiv \theta \square^n \partial_{\mu_1 \dots \mu_l} \theta$

$$\uparrow \Delta^{(0)} = 2\Delta_\theta + 2n + l$$

- DEPENDING AT WHICH ORDER WE ARE, THE INTERMEDIATE OPERATORS ARE DIFFERENT

$$G(\mu, \nu) = \sum_{\Delta, \ell} a_{\Delta, \ell} \mu^{\frac{\Delta-\ell}{2}} \frac{\nu^\ell}{\partial \Delta \ell} g(\mu, \nu)$$

- GENERICALLY:

→ AT LEADING ORDER N^0 ONLY DOUBLE TRACE OPERATORS → $[\Theta\Theta]_{n, \ell} \equiv \Theta \square^n \partial_{\mu_1 \dots \mu_\ell} \Theta$

$$\Delta^{(0)} = 2\Delta_\Theta + 2n + \ell$$

→ AT ORDER N^{-2}

- CORRECTION TO DOUBLE TRACES → $\delta_{n, \ell}^{(1)}$, $a_{n, \ell}^{(1)}$
- NEW OPERATORS → $\tilde{\Theta}$ $\langle \Theta\Theta\tilde{\Theta} \rangle \propto \frac{1}{N}$

→ AT HIGHER ORDERS k

• k -th CORRECTION TO DOUBLE TRACES → $\gamma_{n,e}^{(k)} a_{n,e}^{(k)}$

• HIGHER TRACES $[\theta \theta \dots \theta]$

•

•

• INSERT THESE EXPANSIONS IN THE CONFORMAL BLOCK

DECOMPOSITION:

$$\Delta = 2\Delta_\theta + 2n + \ell$$

$$G_f^{(0)}(u, \sigma) = \sum_{n,e} a_{n,e}^{(0)} u^{\Delta_\theta + n} g_{n,e}(u, \sigma)$$

$$G_f^{(1)}(u, \sigma) = \sum_{n,e} u^{\Delta_\theta + n} \left(a_{n,e}^{(1)} + \frac{1}{2} a_{n,e}^{(0)} \gamma_{n,e}^{(1)} \left(\log u + \frac{\partial}{\partial n} \right) \right) g_{n,e}(u, \sigma)$$

⋮

$$G_f^{(k)}(u, v) = \sum_{n, \ell} \frac{1}{2^k k!} u^{\Delta_0 + n} a_{n, \ell}^{(0)} (\delta_{n, \ell}^{(1)})^k \log^k u g_{n, \ell}(u, v)$$

THIS TALK : UNDERSTAND HOW MUCH INFORMATION (BOTH OF THE OPE DATA AND OF THE AMPLITUDE) WE CAN LEARN FROM THE KNOWLEDGE OF THIS TERM.

WE WILL CONSIDER ONLY ϕ^4 TYPE INTERACTIONS

OTHER MOTIVATIONS :

- TOY MODEL TO UNDERSTAND THE BEHAVIOUR OF HIGHER TRACE OPERATORS OF THE FORM $[000]$ (DIFFERENT FROM KRAVCHUK & MANN)
(FARDELLI FITZPATRICK U)
- EXTEND & COMPLETE THIS APPROACH TO $N=4$ SYM \rightarrow HOW DO TRIPLE TRACES ORGANISE?
(SEE FOR INSTANCE HORRICH, SIMMONS DUFFIN, VIEIRA)
- UNDERSTAND SYSTEMATICS OF LARGE N EXPANSION